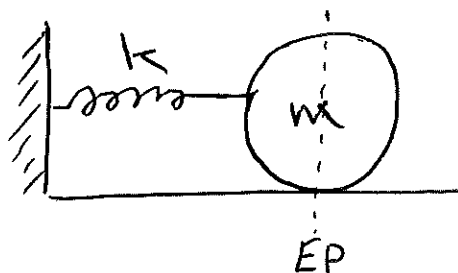


# Classical & Quantum Waves

## Lecture 2-1

### Simple harmonic oscillator (continued)

- Last time we considered SHO of a mass attached to a spring:



- Using Hooke's law,  $F = -kx$ , we arrived at the differential eq.

$$m \frac{d^2x}{dt^2} = -kx \rightarrow \boxed{\frac{d^2x}{dt^2} = -\omega^2 x} \quad \text{with } \omega^2 \equiv \frac{k}{m}$$

- We want to know the solution to this eq.,  $x(t)$

### Displacement, velocity & acceleration in SHO

- Physical intuition (and/or experiments) suggests that a cosine function describes SHO:

Guess that  $x(t) = A \cos\left(2\pi \frac{t}{T}\right)$

$x = A$  at  $t = 0$

Phase angle goes from  $0 \rightarrow 2\pi$  over time  $T$

• Is this a solution to our differential eq.

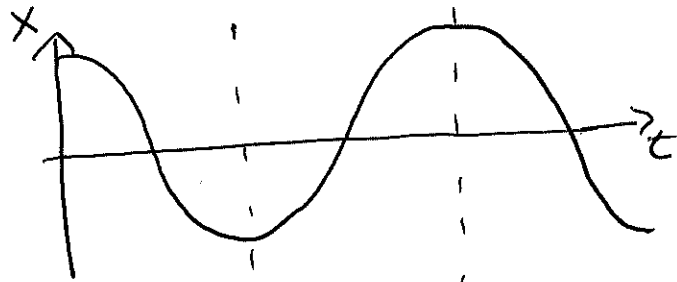
$$\frac{d^2x}{dt^2} = -\omega^2 x \quad ?$$

• Define  $\boxed{\omega \equiv \frac{2\pi}{T}} = 2\pi\nu$

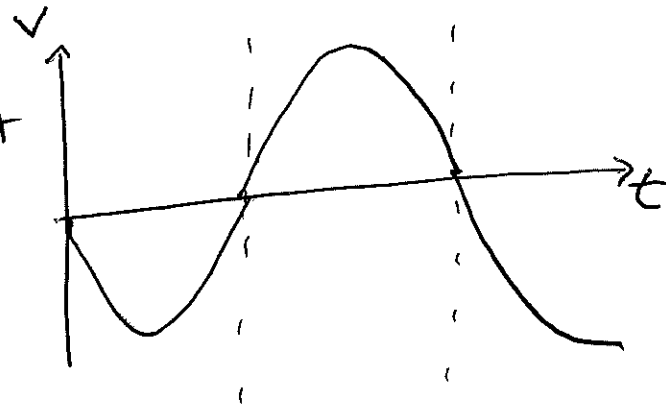
angular frequency: how many  $2\pi$  ~~cycles~~ per sec

~~unit~~: unit:  $[\omega] = \frac{\text{rad}}{\text{s}}$

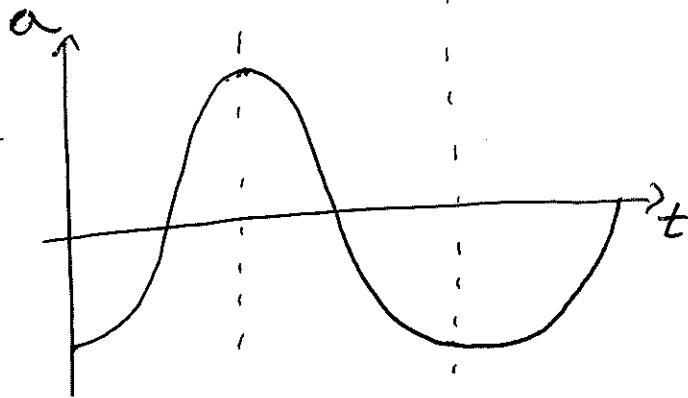
$\rightarrow x = A \cos \omega t$   
displacement



$\rightarrow \frac{dx}{dt} = v = -\omega A \sin \omega t$   
velocity



$\rightarrow \frac{d^2x}{dt^2} = a = -\omega^2 A \cos \omega t$   
 $= -\omega^2 x$   
acceleration

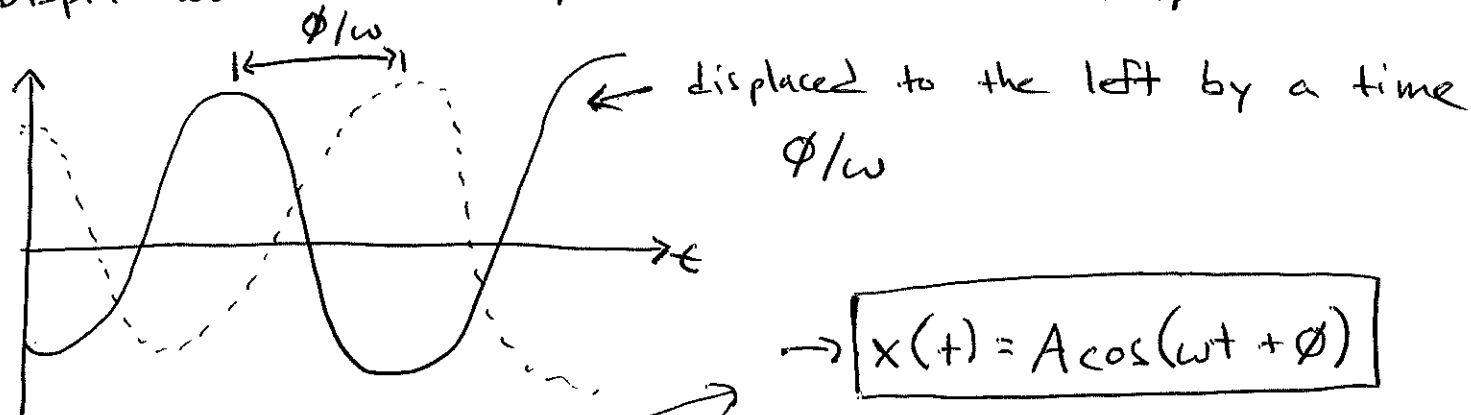


Force & a is max.  
at turning points

What is the general solution for SHO?

2-3

Displacement and velocity at  $t=0$  can be arbitrary!



$$\rightarrow \boxed{x(t) = A \cos(\omega t + \phi)}$$

$\phi$  is called phase angle

• General solution of  $\frac{d^2x}{dt^2} = -\omega^2 x$

• Alternative form:  $\boxed{x(t) = a \cos \omega t + b \sin \omega t}$  (Can show equivalent to  $A \cos(\omega t + \phi)$ )

• Note:  $\omega$  is determined by properties of the oscillator

$\rightarrow$  "natural frequency of oscillation", e.g.  $\omega \equiv \sqrt{\frac{k}{m}}$  for mass on spring

- Oscillation frequency independent of how large initial amplitude

## Energy of a SHO

• Working with the energy is powerful and simpler to calculate

• Mass on a spring has:

• kinetic energy  $K = \frac{1}{2} m v^2$

• potential energy  $U = \int_0^x k x' dx' = \frac{1}{2} k x^2$

(Work to extend spring from  $x'$  to  $x' + dx'$  is  $k x' dx'$ )

• Energy is always conserved, so total energy is

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$$E = K + U = \frac{1}{2}mv^2 + \frac{1}{2}kx^2$$

• Using general solution  $x(t) = A\cos(\omega t + \phi)$ ,  $\frac{dx}{dt} = -\omega A\sin(\omega t + \phi)$

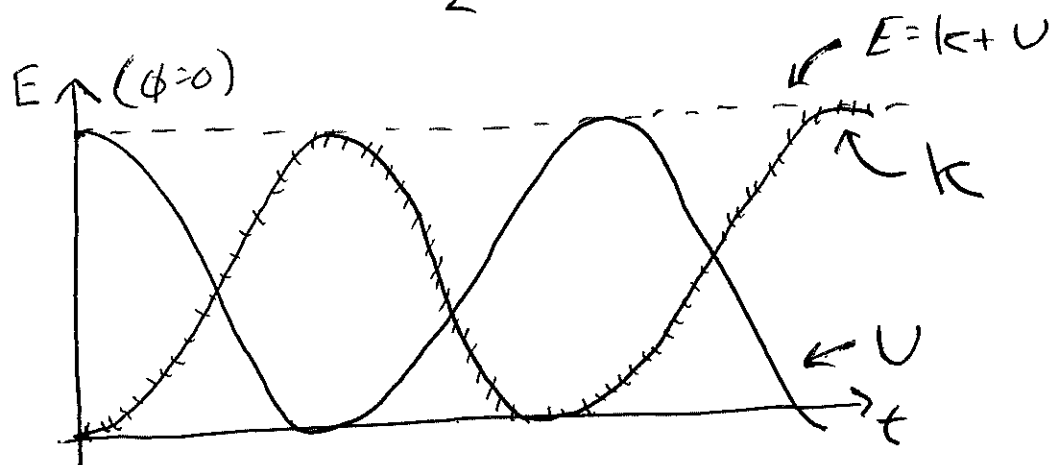
$$U = \frac{1}{2}kx^2 = \frac{1}{2}kA^2\cos^2(\omega t + \phi)$$

$$K = \frac{1}{2}mv^2 = \frac{1}{2}m\omega^2A^2\sin^2(\omega t + \phi) = \frac{1}{2}kA^2\sin^2(\omega t + \phi)$$

$\omega^2 = \frac{k}{m}$

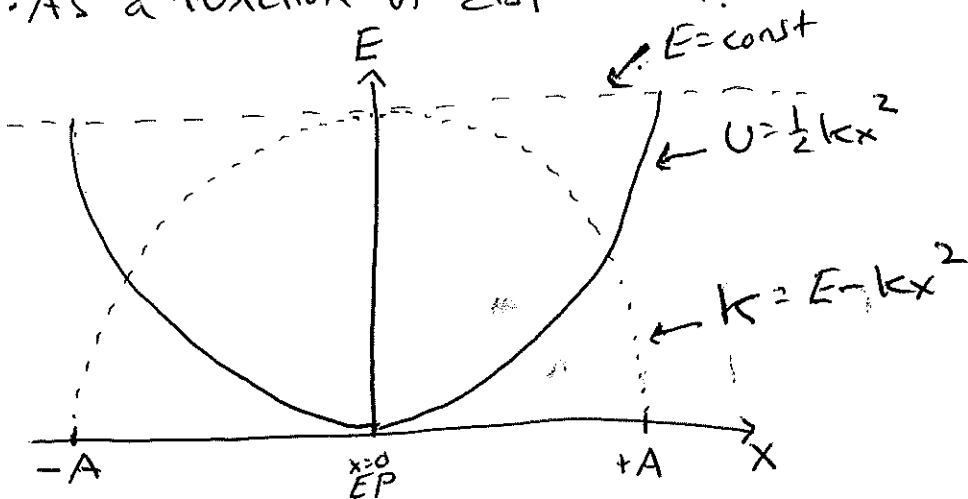
$$\rightarrow E = K + U = \frac{1}{2}kA^2(\underbrace{\sin^2(\omega t + \phi) + \cos^2(\omega t + \phi)}_{=1})$$

$$\rightarrow E = \frac{1}{2}kA^2$$



Energy flows b/t  
K and U!

• As a function of displacement:



• Mass is confined in  
parabolic potential well

↳ Always gives  
SHO!

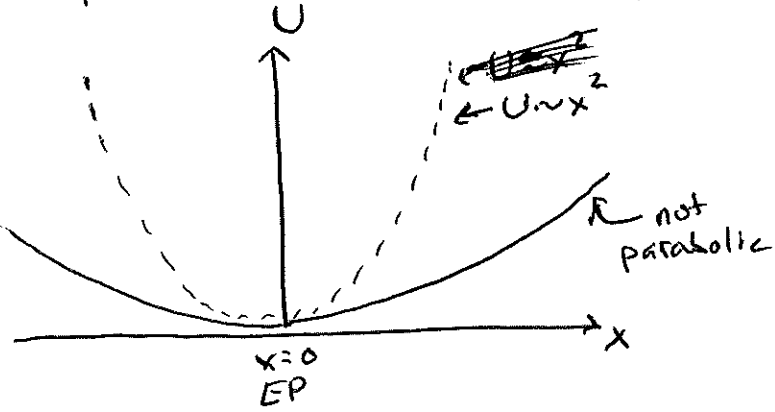
(e.g., marble in a bowl)

What if potential well not parabolic?

Almost all potential wells can be approximated as parabolic close to EP

↳ Why most oscillating systems exhibit SHO when ~~close~~ displacement is small

Explanation: Assume non parabolic potential



Taylor's theorem: A function  $f(x)$  can be expanded around  $x=a$

$$f(x) = f(a) + \frac{(x-a)}{1!} \left( \frac{df}{dx} \right)_{x=a} + \frac{(x-a)^2}{2!} \left( \frac{d^2f}{dx^2} \right)_{x=a} + \dots$$

Expand potential well around EP,  $x=0$ :

$$U(x) = U(0) + x \left( \frac{dU}{dx} \right)_{x=0} + \frac{x^2}{2} \left( \frac{d^2U}{dx^2} \right)_{x=0} + \dots$$

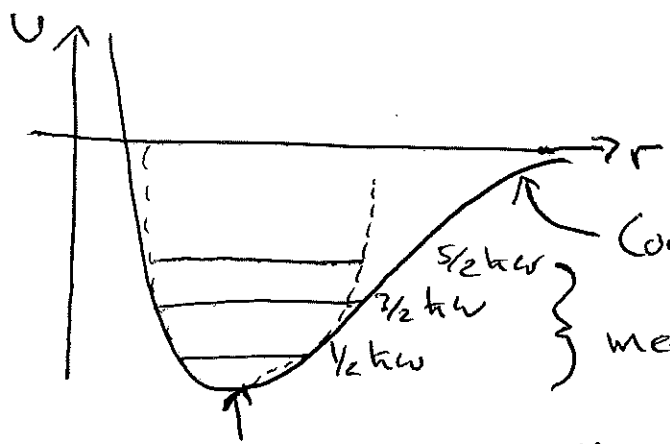
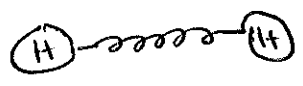
$\uparrow$  Potential energy offset  $\equiv 0$        $\uparrow$   $= 0$  b/c at minimum

For small  $x$ ,  $U(x) \approx \frac{x^2}{2} \left( \frac{d^2U}{dx^2} \right)_{x=0}$   
 $\uparrow$  spring constant  $k$  for mass on spring

Many examples of SHOs around us;

system	Period $T$ [sec]
water in a lake	$10^2 - 10^4$
bridges/buildings	$1 - 10$
clock pendulum	$1$
instruments	$10^{-3} - 10^{-2}$
piezo crystal	$10^{-6}$
molecules	$10^{-15}$

Hydrogen molecule  $H_2$



Coulomb attraction

measurement of vibrational lines yields  $\omega$ !  
(typically infrared)

EP  $r_0 \sim 0.7 \times 10^{-10} \text{ m}$

→ measurement of molecule <sup>optical</sup> spectra gives info on bond strength  ~~$\propto \omega^2$~~   ~~$\propto k/m$~~

↳ get "spring constant" of molecule

(see worked example p. 16-17 King)